

Photoproduction of pentaquark cascades from nucleons

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The cross sections for production of the pentaquark baryon Ξ_5^+ from the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and of Ξ_5^{--} from the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$ are evaluated in a hadronic model that includes their couplings to both $\Sigma \bar{K}$ and $\Sigma \bar{K}^*$. Values for these coupling constants are determined from the empirical $\pi NN(1710)$ and $\rho NN(1710)$ coupling constants by assuming that Ξ_5^+ , Ξ_5^{--} , and $N(1710)$ belong to the same antidecuplet of spin 1/2 and positive parity as the pentaquark Θ^+ . Introducing empirical form factors at strong interaction vertices, we obtain a cross section of about 0.03–0.6 nb for the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and about 0.1–0.6 nb for the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$ at photon energy $E_\gamma = 4.5$ GeV, depending on the value of the coupling constant $g_{K^* \Sigma \Xi_5}$.

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I. INTRODUCTION

The observation of the pentaquark Θ^+ ($uudd\bar{s}$) baryon [1] in nuclear reactions induced by photons [2–5] and kaons [6] has prompted extensive theoretical studies on both its properties [7–22] as well as production [23–30] and decay mechanisms [31,32]. Although most models predict that Θ^+ has spin 1/2 and isospin 0, their predictions on its parity vary widely. While the soliton model gives a positive parity and the lattice QCD study favors a negative parity, the quark model can give either positive or negative parities, depending on whether quarks are correlated or not. The parity of Θ^+ affects both the magnitude of its production cross section in these reactions and the photon asymmetry in photonucleon reactions. With a positive Θ^+ parity, the production cross section is almost an order of magnitude larger than that for a negative parity Θ^+ as a result of a smaller $KN\Theta$ coupling constant in the latter case [25–27]. The positive parity also leads to a large positive photon asymmetry, which becomes negative if the Θ^+ parity is negative [28,30].

There are other pentaquark baryons Ξ_5^+ ($uuss\bar{d}$) and Ξ_5^{--} ($ddss\bar{u}$) in the same antidecuplet as Θ^+ . Although the Ξ_5^{--} has already been observed recently in $p+p$ collisions at center-of-mass energy $\sqrt{s} = 17.2$ GeV by NA49 Collaboration [33], its production in photonucleon reactions has not been studied. In this paper, we shall use the same hadronic model that was introduced for studying Θ^+ production in photonucleon reactions [23–25] to study the production of Ξ_5^+ and Ξ_5^{--} in these reactions.

This paper is organized as follows. In Sec. II, we consider the reactions $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$ for pentaquark cascade production, introduce the interaction Lagrangians needed for evaluating their cross sections, and discuss the coupling constants and form factors at strong interaction vertices. Results on the cross sections for these reactions are then shown and discussed in Sec. III. In Sec. IV, a brief summary is given. Finally, the formula for evaluating the cross section for reactions involving two particles in the initial state and three particles in the final state is derived in the Appendix.

II. PENTAQUARK CASCADE PRODUCTION IN PHOTON-NUCLEON REACTIONS

Possible reactions for pentaquark Ξ_5 production in photon-nucleon reactions near threshold are $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$. To evaluate their cross sections, we use the hadronic model introduced in Refs. [23–25] for studying Θ^+ production in photon-nucleon reactions. This model is a generalization of the SU(3) flavor-invariant Lagrangian for the interactions between octet pseudoscalar mesons and baryons [34] to include the interactions of Θ^+ with nucleons and K as well as K^* . Photon in this model is introduced as $U_{em}(1)$ gauge particle, and the symmetry breaking effects are taken into account phenomenologically by using empirical hadron masses and coupling constants. In the present study, we extend the model to include interactions of the pentaquark Ξ_5 with Σ and K or K^* . Since only three-point interactions appear in the interaction Lagrangians, relevant Feynman diagrams that contribute to these reactions are those shown in Figs. 1 and 2.

A. Interaction Lagrangians

The interaction Lagrangians relevant for evaluating the amplitudes shown in Figs. 1 and 2 are:

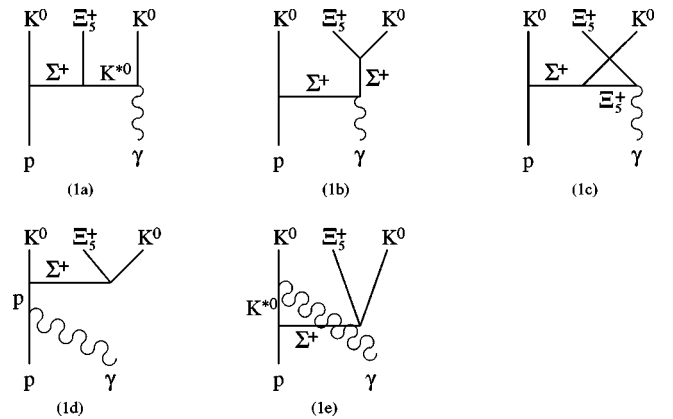


FIG. 1. Diagrams for pentaquark Ξ_5^+ production from the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$.

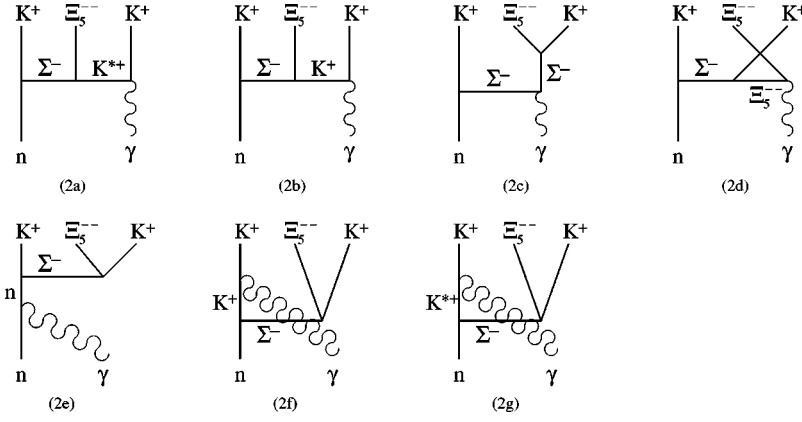


FIG. 2. Diagrams for pentaquark Ξ_5^{--} production from the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$.

$$\mathcal{L}_{KN\Sigma} = ig_{KN\Sigma} \bar{N} \gamma_5 \vec{\Sigma} \cdot \vec{\tau} K + \text{H.c.},$$

$$\mathcal{L}_{K\Sigma\Xi_5} = ig_{K\Sigma\Xi_5} \bar{\Xi}_5 \gamma_5 \vec{\Sigma} \cdot \vec{T} K + \text{H.c.},$$

$$\mathcal{L}_{K^*N\Sigma} = g_{K^*N\Sigma} \bar{N} \gamma_\mu \vec{\Sigma} \cdot \vec{\tau} K^{*\mu} + \text{H.c.},$$

$$\mathcal{L}_{K^*\Sigma\Xi_5} = g_{K^*\Sigma\Xi_5} \bar{\Xi}_5 \gamma_\mu \vec{\Sigma} \cdot \vec{T} K^{*\mu} + \text{H.c.},$$

$$\mathcal{L}_{\gamma NN} = -e\bar{N} \left\{ \gamma_\mu \frac{1+\tau_3}{2} A^\mu - \frac{1}{4m_N} [\kappa_p + \kappa_n + \tau_3(\kappa_p - \kappa_n)] \sigma_{\mu\nu} \partial^\nu A^\mu \right\} N,$$

$$\mathcal{L}_{\gamma\Sigma\Sigma} = -e\bar{\Sigma} \gamma_\mu Q_1 \Sigma A^\mu,$$

$$\mathcal{L}_{\gamma KK} = ie[\bar{K} Q_2 \partial K - (\partial_\mu \bar{K}) Q_2 K] A^\mu,$$

$$\mathcal{L}_{\gamma KK^*} = g_{\gamma KK^*} \epsilon_{\alpha\beta\mu\nu} \partial^\alpha A^\beta [(\partial^\mu \bar{K}^{*\nu}) K + \bar{K} \partial^\mu K^{*\nu}]. \quad (1)$$

In the above, N , Σ , and Ξ_5 denote, respectively, the nucleon isospin doublet, the Σ isospin triplet, and the pentaquark Ξ_5 isospin quartet; K and K^* are the pseudoscalar and vector kaon isospin doublets, respectively; and A_μ denotes the photon. The Pauli matrices are given by $\vec{\tau}$, while \vec{T} is an isospin transition operator represented by a 4×2 matrix with matrix elements given by $\langle \frac{3}{2}m | T_\lambda | \frac{1}{2}n \rangle = \langle \frac{3}{2}m | 1\lambda \frac{1}{2}n \rangle$, where $\lambda = 0, \pm 1$. The operators $Q_1 = \text{diag}(1, 0, -1)$ and $Q_2 = \text{diag}(1, 0)$ are diagonal charge operators and $\epsilon_{\alpha\beta\mu\nu}$ denotes the antisymmetric tensor with the usual convention $\epsilon_{0123} = 1$.

For the coupling constants $g_{KN\Sigma}$ and $g_{K^*N\Sigma}$, their values can be related to $g_{\pi NN}$ and $g_{\rho NN}$, which have empirical values of $g_{\pi NN} = 13.5$ [35] and $g_{\rho NN} = 3.25$ [36], by SU(3) symmetry, i.e., $g_{KN\Sigma} = (1 - 2\alpha)g_{\pi NN} = -3.78$ and $g_{K^*N\Sigma} = g_{\rho NN} = 3.25$ if we use the ratio $\alpha = D/(D+F) = 0.64$ [37] for the D - and F -type interaction Lagrangians between pseudoscalar mesons and baryons. Similarly, we have $g_{K\Sigma\Xi_5} = g_{KN\Theta}/\sqrt{2} = 2.16$ and $g_{K^*\Sigma\Xi_5} = g_{K^*N\Theta}/\sqrt{2} = 1.27$ [38], where the values for $g_{KN\Theta} = 3.06$ and $g_{K^*N\Theta} = 1.8$ have been determined in Ref. [25]

from the empirical $\pi NN(1710)$ and $\rho NN(1710)$ coupling constants using the SU(3) symmetry. Since the sign of $g_{K^*N\Theta}$ relative to $g_{KN\Theta}$ and $g_{KN\Sigma}$ is not fixed by the SU(3) symmetry, we shall consider both signs for the coupling constant $g_{K^*\Sigma\Xi_5} = \pm 1.27$ as well as the case of $g_{K^*\Sigma\Xi_5} = 0$.

For photon coupling to nucleon, we include also its interaction with the anomalous magnetic moment of nucleons with empirical values of $\kappa_p = 1.79$ and $\kappa_n = -1.91$. Since the anomalous magnetic moment of Θ^+ is not known, we neglect its coupling to photon.

The coupling constant $g_{\gamma KK^*}$ denotes the photon anomalous parity interaction with kaons and has the dimension of inverse of energy. Its value is $g_{\gamma K^0 K^{*0}} = 0.388 \text{ GeV}^{-1}$ and $g_{\gamma K^\pm K^{*\pm}} = 0.254 \text{ GeV}^{-1}$ [24] using the decay width $\Gamma_{K^{*0} \rightarrow K^0 \gamma} = 0.117 \text{ MeV}$ of K^{*0} and $\Gamma_{K^{*\pm} \rightarrow K^\pm \gamma} = 0.05 \text{ MeV}$ of $K^{*\pm}$ to kaon and photon [39]. Although the sign of $g_{\gamma KK^*}$ relative to other coupling constants in the interaction Lagrangians is not known either, it is not relevant for our study as both constructive and destructive interferences among the diagrams in Fig. 1 and 2 are automatically taken into account by using different signs for the coupling constant $g_{K^*\Sigma\Xi_5}$.

B. Amplitudes for the reactions

$\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^{--}$

With the above interaction Lagrangians, we can write the amplitudes for the five diagrams in Fig. 1 as

$$\begin{aligned} \mathcal{M}_{1a} &= 2ig_{\gamma K^0 K^{*0}} g_{K^*\Sigma\Xi_5} g_{KN\Sigma} \frac{1}{[(p_2 - p_5)^2 - m_{K^*}^2](t - m_\Sigma^2)} \\ &\quad \times \epsilon_{\alpha\beta\mu\nu} p_2^\alpha p_5^\beta \epsilon^\mu \bar{\Xi}_5(p_4) \gamma^\nu (\not{p}_1 - \not{p}_3 + m_\Sigma) \gamma_5 p(p_1), \\ \mathcal{M}_{1b} &= -2eg_{K\Sigma\Xi_5} g_{KN\Sigma} \frac{1}{(s_1 - m_\Sigma^2)(t - m_\Sigma^2)} \\ &\quad \times \bar{\Xi}_5(p_4) (\not{p}_4 + \not{p}_5 - m_\Sigma) \not{\epsilon} (\not{p}_1 - \not{p}_3 - m_\Sigma) p(p_1), \\ \mathcal{M}_{1c} &= -2eg_{K\Sigma\Xi_5} g_{KN\Sigma} \frac{1}{[(p_2 - p_4)^2 - m_{\Xi_5}^2](t - m_\Sigma^2)} \\ &\quad \times \bar{\Xi}_5(p_4) \not{\epsilon} (\not{p}_4 - \not{p}_2 + m_{\Xi_5}) (\not{p}_1 - \not{p}_3 - m_\Sigma) p(p_1), \end{aligned}$$

$$\begin{aligned}
\mathcal{M}_{1d} &= 2eg_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{(s_1-m_\Sigma^2)(s-m_N^2)} \\
&\quad \times \bar{\Xi}(p_4)(\not{p}_4+\not{p}_5-m_\Sigma)(\not{p}_1+\not{p}_2+m_N) \\
&\quad \times \left(\not{\epsilon}-\frac{\kappa_p}{2m_N}\not{\epsilon}\not{p}_2\right)p(p_1), \\
\mathcal{M}_{1e} &= 2ig_{K\Sigma\Xi_5}g_{K^*N\Sigma}g_{\gamma K^0 K^{*0}}\frac{1}{(s_1-m_\Sigma^2)[(p_2-p_3)^2-m_{K^*}^2]} \\
&\quad \times \epsilon_{\alpha\beta\mu\nu}p_2^\alpha p_3^\beta \epsilon^\mu \bar{\Xi}_5(p_4)\gamma_5(\not{p}_4+\not{p}_5+m_\Sigma)\gamma^\nu p(p_1),
\end{aligned} \tag{2}$$

and those for the seven diagrams in Fig. 2 as

$$\begin{aligned}
\mathcal{M}_{2a} &= 2ig_{\gamma K^+ K^{*+}}g_{K^*N\Sigma}g_{K^*N\Sigma}\frac{1}{[(p_2-p_5)^2-m_{K^*}^2](t-m_\Sigma^2)} \\
&\quad \times \epsilon_{\alpha\beta\mu\nu}p_2^\alpha p_5^\beta \epsilon^\mu \bar{\Xi}_5(p_4)\gamma^\nu(\not{p}_1-\not{p}_3+m_\Sigma)\gamma_5 n(p_1), \\
\mathcal{M}_{2b} &= -4eg_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{[(p_2-p_5)^2-m_K^2](t-m_\Sigma^2)}\epsilon p_5\bar{\Xi}_5(p_4) \\
&\quad \times (\not{p}_1-\not{p}_3-m_\Sigma)n(p_1), \\
\mathcal{M}_{2c} &= 2eg_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{(s_1-m_\Sigma^2)(t-m_\Sigma^2)} \\
&\quad \times \bar{\Xi}_5(p_4)(\not{p}_4+\not{p}_5-m_\Sigma)\not{\epsilon}(\not{p}_1-\not{p}_3-m_\Sigma)n(p_1), \\
\mathcal{M}_{2d} &= 4eg_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{[(p_2-p_4)^2-m_{\Xi_5}^2](t-m_\Sigma^2)} \\
&\quad \times \bar{\Xi}_5(p_4)\not{\epsilon}(\not{p}_4-\not{p}_2+m_{\Xi_5})(\not{p}_1-\not{p}_3-m_\Sigma)n(p_1), \\
\mathcal{M}_{2e} &= -\frac{e\kappa_n}{m_N}g_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{(s_1-m_\Sigma^2)(s-m_N^2)}\bar{\Xi}_5(p_4)(\not{p}_4+\not{p}_5 \\
&\quad -m_\Sigma)(\not{p}_1+\not{p}_2+m_N)\not{\epsilon}\not{p}_2 n(p_1), \\
\mathcal{M}_{2f} &= -4eg_{K\Sigma\Xi_5}g_{K^*N\Sigma}\frac{1}{(s_1-m_\Sigma^2)[(p_2-p_3)^2-m_K^2]}\epsilon p_3 \\
&\quad \times \bar{\Xi}_5(p_4)(\not{p}_4+\not{p}_5-m_\Sigma)n(p_1), \\
\mathcal{M}_{2g} &= 2ig_{K\Sigma\Xi_5}g_{K^*N\Sigma}g_{\gamma K^+ K^{*+}}\frac{1}{(s_1-m_\Sigma^2)[(p_2-p_3)^2-m_{K^*}^2]} \\
&\quad \times \epsilon_{\alpha\beta\mu\nu}p_2^\alpha p_3^\beta \epsilon^\mu \bar{\Xi}_5(p_4)\gamma^5(\not{p}_4+\not{p}_5+m_\Sigma)\gamma^\nu n(p_1).
\end{aligned} \tag{3}$$

In the above, p_1 and p_2 denote, respectively, the momenta of nucleon and photon in the initial states; p_3 , p_4 , and p_5 are, respectively, those of the kaon on the left, the cascade, and the kaon on the right in the final states of the Feynman dia-

grams in Figs. 1 and 2. We have also introduced the photon polarization vector ϵ^μ and the definitions $s=(p_1+p_2)^2$, $t=(p_1-p_3)^2$, and $s_1=(p_4+p_5)^2$.

C. Form factors

To take into account the effects due to hadron internal structure, a form factor is introduced to each amplitude, and it is taken to have the following form [40]:

$$F(m_x, m_y) = \left[\frac{\Lambda^4}{\Lambda^4 + (q_x^2 - m_x^2)^2} \right] \left[\frac{\Lambda^4}{\Lambda^4 + (q_y^2 - m_y^2)^2} \right], \tag{4}$$

where q_x and q_y are four-momenta of the intermediate off-shell particles with masses m_x and m_y in each diagram, and Λ is the cutoff parameter that characterizes the off-shell momentum above which hadron internal structure becomes important. Since the amplitudes for diagrams (1a) and (1e) in Fig. 1 for the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ are individually gauge invariant, including form factors does not affect their gauge invariance. This is different for the amplitudes for diagrams (1b)–(1d) as only their sum is gauge invariant. Using different form factors for these amplitudes thus leads to a violation of the gauge invariance. To recover the gauge invariance for the total amplitude, one can either drop the gauge violating terms in the amplitude [41] or use an averaged form factor [42] or other combinations of form factors [43] for these terms. For simplicity, we use here a common form factor, which is given by the average of the original form factors for diagrams (1b)–(1d), for all these three diagrams. A similar consideration is applied to the diagrams in Fig. 2 for the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^-$. In the following, we summarize the form factors used in the present study:

$$\begin{aligned}
F_{(1a)} &= F_{(2a)} = F(m_\Sigma, m_{K^*}), \\
F_{(1b)} &= F_{(1c)} = F_{(1d)} = \frac{1}{3}[F(m_\Sigma, m_\Sigma) + F(m_\Sigma, m_{\Xi_5}) \\
&\quad + F(m_\Sigma, m_N)], \\
F_{(1e)} &= F_{(2g)} = F(m_{K^*}, m_\Sigma), \\
F_{(2b)} &= F_{(2c)} = F_{(2d)} = F_{(2f)} = \frac{1}{4}[F(m_\Sigma, m_K) + F(m_\Sigma, m_\Sigma) \\
&\quad + F(m_\Sigma, m_{\Xi_5}) + F(m_K, m_\Sigma)], \\
F_{(2e)} &= F(m_\Sigma, m_N),
\end{aligned} \tag{5}$$

where the subscripts refer to the diagrams in Figs. 1 and 2.

The value of the cutoff parameter is taken to be $\Lambda = 1.7$ GeV, which is determined from fitting the measured cross section for charmed hadron production with three-body final states from photon-proton reactions at center-of-mass energy of 6 GeV [44] using a similar hadronic model based on SU(4) flavor-invariant Lagrangians with empirical hadron masses and coupling constants. This cutoff parameter is larger than the value of 1.2 GeV used in our previous study of the reaction $\gamma p \rightarrow K^{*0} \Theta^+$ [25], where the form factor was taken to be $F = \Lambda^4/(\Lambda^4 + m_K^4)$ and the cutoff parameter was determined from fitting the cross section for charmed hadron

production with two-body final states in photon-proton reactions.

D. Cross sections for $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^-$

As shown in the Appendix, the total cross section for these reactions can be expressed as

$$\sigma_{\gamma N \rightarrow KK\Xi_5} = \frac{1}{(2\pi)^4} \frac{1}{128sp^2} \int_{s_-}^{s_+} ds_1 \int_{t_-}^{t_+} dt \frac{k}{2\sqrt{s_1}} \times \int_0^\pi \sin \theta d\theta \int_0^\pi d\phi \left| \sum_i \mathcal{M}_i(\gamma N \rightarrow KK\Xi_5) \right|^2, \quad (6)$$

where s , t , and s_1 were defined after Eq. (3). The limits of integrations are

$$s_- = (m_{\Xi_5} + m_K)^2, \quad s_+ = (s^{1/2} - m_K)^2,$$

$$t_{\pm} = m_N^2 + m_K^2 - \frac{(m_K^2 + s - s_1)\sqrt{m_N^2 + p^2}}{\sqrt{s}} \pm 2p \sqrt{\frac{(m_K^2 + s - s_1)^2}{4s} - m_K^2}, \quad (7)$$

where p is the momentum of the nucleon in the total center-of-mass frame. We have also introduced by \mathbf{k} the momentum of Ξ_5 in the center-of-mass frame of Ξ_5 and the kaon on the right of the Feynman diagram, with polar angles θ and ϕ , i.e., the angles θ_4 and ϕ_4 shown in Fig. 5 of the Appendix.

III. RESULTS

We first show in the upper panels of Figs. 3 and 4 the cross sections for the reactions $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ as functions of photon energy in the laboratory frame without including form factors at the strong interaction vertices. Both cross sections are seen to increase rapidly with increasing photon energy. Since the cross sections for the case of $g_{K^*\Sigma\Xi_5}=0$ are between those with $g_{K^*\Sigma\Xi_5}=1.27$ and $g_{K^*\Sigma\Xi_5}=-1.27$, the contribution from diagrams involving the coupling $g_{K^*\Sigma\Xi_5}$ is less important than those involving the coupling $g_{K\Sigma\Xi_5}$. In the lower panels, the cross sections with form factors are shown, and they are much smaller than those without form factors. Depending on the value of the coupling constant $g_{K^*\Sigma\Xi_5}$, the cross section for the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ has values of 0.03–0.6 nb at $E_\gamma=4.5$ GeV, while that for the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ has values of 0.1–0.6 nb. We note that these values are significantly smaller than those predicted for Θ^+ production in photon-nucleon reactions using the same hadronic model [23–25].

IV. SUMMARY

Using a hadronic model that includes the coupling of pentaquark Ξ_5 to usual Σ and K or K^* , we have evaluated the cross sections for their production in the reactions $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ by assuming that Ξ_5 has spin

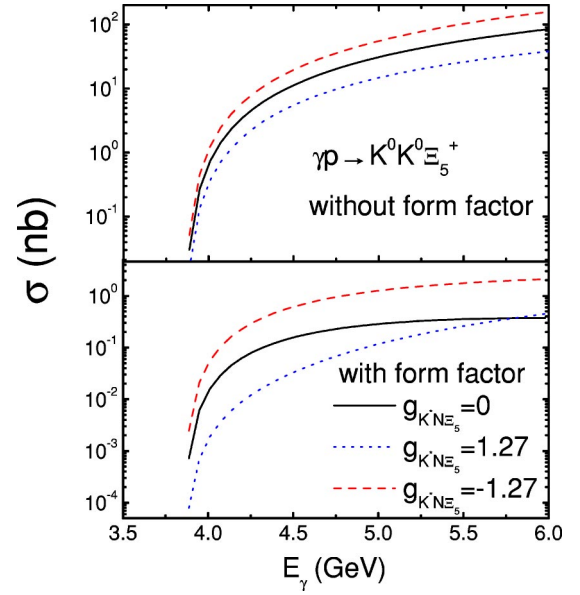


FIG. 3. (Color online) Total cross section for Ξ_5^+ production from the reaction $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ as a function of photon energy and for the coupling constant $g_{K^*\Sigma\Xi_5}=1.27$ (dotted curve), 0 (solid curve), and -1.27 (dashed curve). Upper and lower panels are for the cases without and with form factors, respectively.

1/2 and positive parity. Using coupling constants related to those for pentaquark Θ^+ couplings to N and K or K^* , and also including form factors at the strong interaction vertices with empirical cutoff parameters, these cross sections are found in the range of 0.03–0.6 nb for $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and 0.1–0.6 nb for $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ at photon energy $E_\gamma=4.5$ GeV. Since the coupling constant $g_{KN\Theta}$ is much smaller for a negative parity

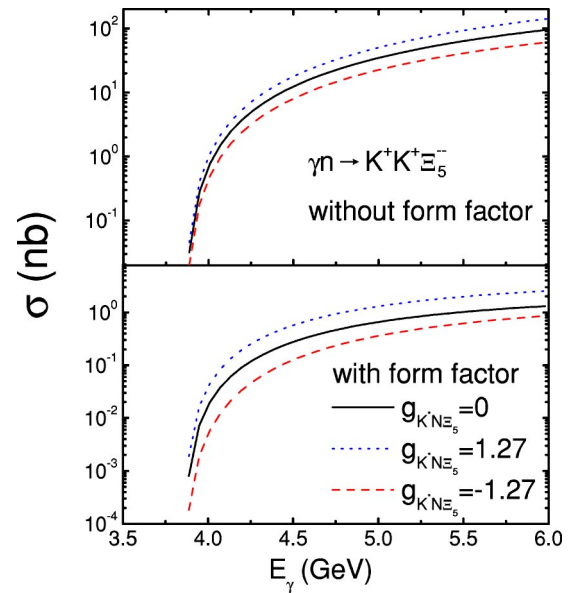


FIG. 4. (Color online) Total cross section for Ξ_5^- production from the reaction $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ as a function of photon energy and for the coupling constant $g_{K^*\Sigma\Xi_5}=1.27$ (dotted curve), 0 (solid curve), and -1.27 (dashed curve). Upper and lower panels are for the case without and with form factors, respectively.

Θ^+ than for a positive parity one, the coupling constant $g_{K\Sigma\Xi_5}$ also becomes smaller if Ξ_5 has a negative parity. As a result, the cross sections for producing negative pentaquark Ξ_5^+ and Ξ_5^- are much smaller than for positive ones as in the case of Θ^+ production. Although the cross sections for photoproduction of pentaquark Ξ_5 are an order of magnitude smaller than those for producing the Θ^+ in these reactions, study of Ξ_5 production will be useful for understanding the properties of pentaquark baryons.

ACKNOWLEDGMENTS

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APPENDIX

In this appendix, we describe the derivation of the cross section formula shown in Eq. (6) for the reactions $\gamma p \rightarrow K^0 K^0 \Xi_5^+$ and $\gamma n \rightarrow K^+ K^+ \Xi_5^-$ involving two particles in the initial state and three particles in the final state, i.e., $1+2 \rightarrow 3+4+5$. It essentially follows the method given in Ref. [45].

In the center-of-mass frame of particles 1 and 2, the cross section for such a reaction generally reads as

$$\begin{aligned} \sigma_{1+2 \rightarrow 3+4+5} = & \frac{1}{(2\pi)^5} \frac{1}{4s^{1/2}p} \int \int d^4 p_3 d^4 p_4 \delta(p_3^2 - m_3^2) \\ & \times \delta^{(4)}(p_\Delta - p_1 + p_3) \int \int d^4 p_4 d^4 p_5 \delta(p_4^2 - m_4^2) \\ & \times \delta(p_5^2 - m_5^2) \delta^{(4)}(q - p_4 - p_5) \\ & \times |\mathcal{M}(p_1 + p_2 \rightarrow p_3 + p_4 + p_5)|^2, \end{aligned} \quad (\text{A1})$$

with $s = (p_1 + p_2)^2$, $q = p_1 + p_2 - p_3$, and p denoting the momentum of initial particles in their center-of-mass frame.

To evaluate the phase-space integral in Eq. (A1), we choose the coordinate system in the center-of-mass frame of particles 4 and 5 as shown in Fig. 5, in which the z axis is along the three-momentum \mathbf{p}_2 of particle 2, while \mathbf{p}_1 and \mathbf{p}_3 are in the x - z plane. In terms of the angles shown in the figure, the four-momenta of these particles can be expressed as

$$p_1 = (E_1, |\mathbf{p}_3| \sin \theta_3, 0, |\mathbf{p}_3| \cos \theta_3 - |\mathbf{p}_2|),$$

$$p_2 = (E_2, 0, 0, |\mathbf{p}_2|),$$

$$p_3 = (E_3, |\mathbf{p}_3| \sin \theta_3, 0, |\mathbf{p}_3| \cos \theta_3),$$

$$p_4 = (E_4, |\mathbf{p}_4| \sin \theta_4 \sin \phi_4, |\mathbf{p}_4| \sin \theta_4 \cos \phi_4, |\mathbf{p}_4| \cos \theta_4),$$

$$\begin{aligned} p_5 = & (E_5, -|\mathbf{p}_4| \sin \theta_4 \sin \phi_4, -|\mathbf{p}_4| \sin \theta_4 \cos \phi_4, \\ & -|\mathbf{p}_4| \cos \theta_4). \end{aligned} \quad (\text{A2})$$

From four-momentum conservation and on-shell constraints, the following identities can be derived:

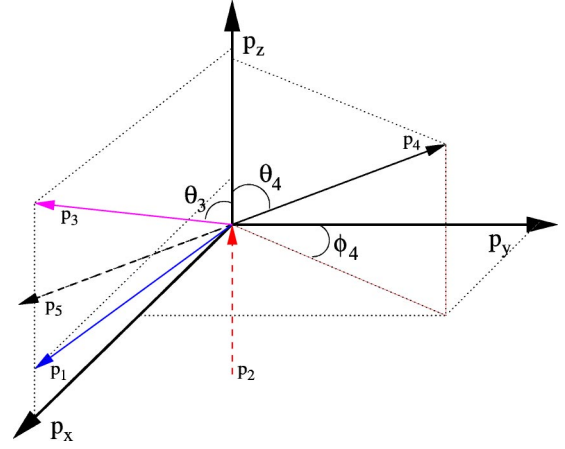


FIG. 5. (Color online) Coordinate system used for the three-momenta of the particles in the reaction $1+2 \rightarrow 3+4+5$.

$$E_1 = \frac{t + s - m_2^2 - m_3^2}{2\sqrt{s_1}}, \quad E_2 = \frac{s_1 - t + m_2^2}{2\sqrt{s_1}}, \quad E_3 = \frac{s - s_1 - m_3^2}{2\sqrt{s_1}},$$

$$E_4 = \frac{s_1 + m_4^2 - m_5^2}{2\sqrt{s_1}}, \quad E_5 = \frac{s_1 + m_5^2 - m_4^2}{2\sqrt{s_1}},$$

$$\cos \theta_3 = \frac{m_1^2 - E_1^2 + |\mathbf{p}_3|^2 + |\mathbf{p}_2|^2}{2|\mathbf{p}_3||\mathbf{p}_2|},$$

$$|\mathbf{p}_2| = \sqrt{E_2^2 - m_2^2}, \quad |\mathbf{p}_3| = \sqrt{E_3^2 - m_3^2}, \quad |\mathbf{p}_4| = \sqrt{E_4^2 - m_4^2}, \quad (\text{A3})$$

where $t = (p_1 - p_3)^2$ and $s_1 = (p_4 + p_5)^2$. The two-body phase-space integral can be rewritten as

$$\begin{aligned} & \int d^4 p_4 d^4 p_5 \delta(p_4^2 - m_4^2) \delta(p_5^2 - m_5^2) \delta^{(4)}(q - p_4 - p_5) \\ & = \frac{k}{2\sqrt{s_1}} \int_0^\pi \sin \theta_4 d\theta_4 \int_0^\pi d\phi_4, \end{aligned} \quad (\text{A4})$$

where k is the momentum of particle 4 or 5 in their center-of-mass frame.

After integrating over the four-momentum p_3 , we have

$$\begin{aligned} \sigma_{1+2 \rightarrow 3+4+5} = & \frac{1}{(2\pi)^5} \frac{1}{4s^{1/2}p} \int d^4 p_\Delta \delta((p_1 - p_\Delta)^2 - m_3^2) \\ & \times \frac{k}{2\sqrt{s_1}} \int_0^\pi \sin \theta_4 d\theta_4 \int_0^\pi d\phi_4 |\mathcal{M}(p_1 + p_2 \rightarrow p_3 \\ & + p_4 + p_5)|^2. \end{aligned} \quad (\text{A5})$$

Denoting $p_\Delta = (E_\Delta, \mathbf{p}_\Delta)$, and then using the identities

$$s_1 = m_3^2 - m_1^2 + m_2^2 + 2\sqrt{s}E_\Delta,$$

$$t \equiv p_\Delta^2 = m_3^2 - m_1^2 + 2\sqrt{m_1^2 + p^2}E_\Delta - 2pp_{\Delta z} \quad (\text{A6})$$

in the center-of-mass frame of particles 1 and 2, we can rewrite

$$\int d^4p_\Delta \delta((p_1 - p_\Delta)^2 - m_3^2) = \frac{\pi}{4ps^{1/2}} \int ds_1 \int dt. \quad (\text{A7})$$

We then have

$$\begin{aligned} \sigma_{1+2 \rightarrow 3+4+5} &= \frac{1}{(2\pi)^4} \frac{1}{32sp^2} \int_{s_-}^{s_+} ds_1 \int_{t_-}^{t_+} dt \\ &\times \frac{k}{2\sqrt{s_1}} \int_0^\pi \sin \theta_4 d\theta_4 \int_0^\pi d\phi_4 |\mathcal{M}(p_1 + p_2 \rightarrow p_3 \\ &+ p_4 + p_5)|^2, \end{aligned} \quad (\text{A8})$$

where

$$\begin{aligned} s_- &= (m_4 + m_5)^2, s_+ = (s^{1/2} - m_3)^2, \\ t_\pm &= m_1^2 + m_3^2 \\ &- \frac{(m_3^2 + s - s_1)\sqrt{m_1^2 + p^2}}{\sqrt{s}} \pm 2p \sqrt{\frac{(m_3^2 + s - s_1)^2}{4s} - m_3^2}. \end{aligned} \quad (\text{A9})$$

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